MEASUREMENT OF INDUCTANCE BY ANDERSON'S METHOD, USING ALTERNATING CURRENTS AND A VIBRATION GALVANOMETER.

By Edward B. Rosa and Frederick W. Grover.

1. HISTORY OF THE METHOD.

Several modifications of Maxwell's method[a] of comparing an inductance with a capacity have been proposed in order to obviate the double adjustment of resistances necessary in that method. Maxwell showed
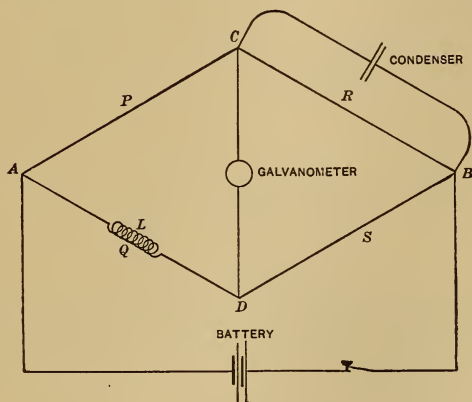

FIG. 1.—Maxwell's method.

that if (1) the bridge is balanced for steady currents and at the same time (2) the resistances are so chosen that there is no deflection of the galvanometer when the battery current is suddenly closed or broken, then

$$L = CRQ = CPS \qquad (1)$$

where L is the inductance in the arm $A\ D$, the resistance of which is

---

[a] Electricity and Magnetism, § 778.

Q, C is the value of the capacity in parallel with R, and P, R, and S are noninductive resistances.

In order to satisfy both of these conditions two of the arms of the bridge must be varied simultaneously, so that the balance for steady currents may be maintained while the balance for transient currents is sought. This is generally a tedious process, although by means of a small variable inductance in Q, in addition to the inductance to be measured, and a multiple valued condenser the process might be considerably accelerated.

In 1891 Professor Anderson proposed[a] an important modification of Maxwell's method, which consisted in joining the condenser to a point E, separated from C by a variable resistance r. The bridge being balanced for steady currents by varying any one of the four arms of the bridge, the balance for transient currents is then made by
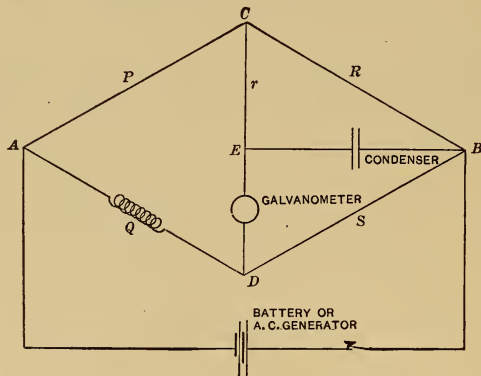

FIG. 2.—Anderson's method.

varying r, which does not disturb the balance of the bridge for steady currents. This change, which rendered the two adjustments independent, removed at once a most serious difficulty and made the method thoroughly practicable.

Anderson's demonstration for the case of transient currents gives for the value of the inductance (changing the letters to correspond to fig. 2)

$$L = C\,[r\,(Q+S)+PS] \qquad (2)$$

If $r=0$, $L=CPS$, as in Maxwell's method.

In the use of Anderson's method r may be small, so that CPS is the principal part of the expression for the inductance, or it may be larger, and the first term, $Cr\,(Q+S)$, represents the larger part of L. Thus a considerable range of values of inductance may be measured

a Phil. Mag., 31, p. 329, 1891.

without changing the arms of the bridge or the capacity of the condenser.

Stroud and Oates[a] have proposed another modification of Maxwell's method, which they have used with much success in measuring inductances. Instead of employing an interrupted current from a battery, as Anderson had done, they used an alternating current and an alternating-current galvanometer, the latter being essentially a d'Arsonval galvanometer, with the field magnet laminated and strongly excited by an alternating current from the generator. The galvanometer was thus made very sensitive, and to increase the sensitiveness still further the resistance r was placed outside the bridge, as shown in Fig. 3. It will be seen that this arrangement differs from Maxwell's only in separating the point B from the terminal of the condenser by the
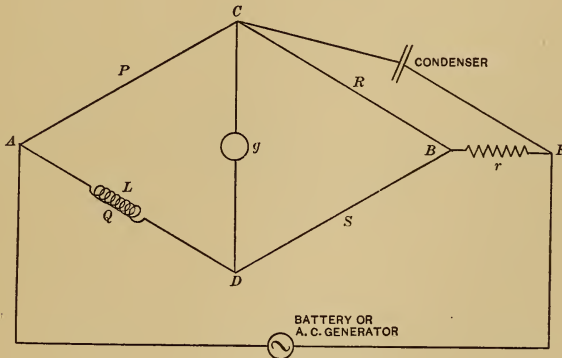

FIG. 3.—Stroud's method.

auxiliary adjustable resistance r, which in Anderson's method is in the galvanometer circuit between C and D. As the resistance r is sometimes several hundred ohms, it reduces the sensibility when in the galvanometer circuit, whereas in the arrangement of Fig. 3 the electromotive force can be increased if r is large, and so keep the same current in the bridge as when r is small, and thus maintain the sensibility.

The expression for the inductance L in Stroud's method (changing the letters to correspond with Fig. 3) is

$$L = C\,[r\,(Q+P)+PS], \tag{3}$$

which closely resembles the formula for Anderson's method, but differs in having $Q+P$ in the first term instead of $Q+S$.

---

[a] Phil. Mag., **6**, p. 707, 1903.

Professsor Fleming has pointed out that Stroud's arrangement may be regarded as conjugate to Anderson's, the galvanometer and source of current being interchanged, Fig. 4. In this case the formula is exactly the same as for Anderson's method. If, however, Fig. 4 be rearranged so as to agree with Fig. 3, it will be found that the arms P and S are interchanged, and consequently that these letters must be interchanged in the formula for L. This changes equation (2) into equation (3).

Fleming and Clinton have employed Anderson's method for the measurement of small inductances, using a battery and a rotating commutator and galvanometer,[a] and later Fleming employed an interrupted current, produced by a vibrating armature, and a telephone.[b]
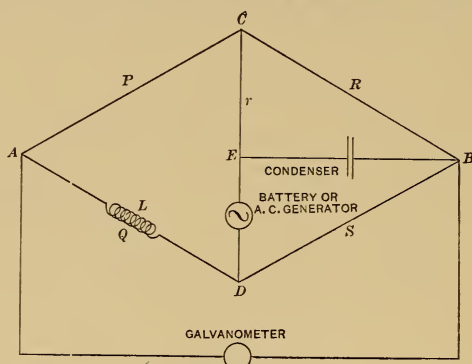

FIG. 4.—Showing Stroud's method as conjugate to Anderson's.

During the past two years we have employed Anderson's method for the measurement of both large and small inductances, using (1) a battery as a source of current and a d'Arsonval galvanometer, with a rotating commutator to interrupt and reverse simultaneously the current and galvanometer terminals; or (2), what has proved more satisfactory, an alternating current and a vibration galvanometer, the latter being tuned to the frequency of the current furnished by the generator.

2. ADVANTAGES OF THE METHOD.

We have found the method rapid and convenient in practice and the vibration galvanometer sufficiently sensitive to permit very accurate settings. As compared with other methods of accurately measuring inductance, it possesses striking advantages, some of which will here be specifically mentioned.

---

[a] Phil. Mag., 5, p. 493; 1903.   [b] Phil. Mag., 7, p. 586; 1904.

(*a*) All methods of measuring inductances without the use of a condenser (or other known inductance) require an accurate knowledge of the frequency of the alternating current employed. It is not difficult to determine accurately the mean frequency of an alternating current, even when the generator is inaccessible, as a counter may be employed to record on a chronograph the number of revolutions in a given time; moreover, the speed of the generator may be maintained sufficiently constant to enable good settings to be made. But to hold the speed steady enough to make settings of a high order of accuracy is difficult and requires an assistant to control the speed. With Anderson's method, even with a tuned galvanometer, *slight* changes of frequency are not detrimental, and hence the labor of taking the observations is greatly reduced.

(*b*) The inductance is determined in terms of a capacity, in addition to several resistances, which are also required in other methods of measuring inductance. A capacity can be measured by Maxwell's bridge method, using a commutator, with very great exactness, provided care is taken in choosing the resistances of the arms of the bridge,[a] and also provided the temperature of the condenser is taken and a temperature correction subsequently applied whenever necessary. The capacity of a condenser is not the same for slow charges as for rapid charges, and hence, if Anderson's method is used for transient currents, the capacity employed in the formula should correspond to the conditions of the experiment. As the successive makes and breaks of the current are likely to be irregular, the result would be that the effective capacity would vary slightly in successive trials, even with the best mica condensers. On the other hand, using an interrupted or alternating current of constant frequency, the capacity is uniform and definite, and if it is measured at the same frequency there is no uncertainty as to its value. In our experiments we employ an eight-pole generator, giving four complete cycles in each revolution. To this generator is joined the commutator which is employed in charging and discharging the condenser when measuring its capacity, the commutator having four segments, and hence charging and discharging the condenser four times in each revolution. Thus the frequency of charge and discharge of the condenser may be made exactly the same in use as when its capacity is measured. The change of capacity of a condenser with the frequency is very slight, but in measurements of the highest accuracy it is well to eliminate the slight uncertainty due to change of frequency.

---

[a] Bulletin of the Bureau of Standards, No. 2, 1905.

(c) The formula for calculating the inductance is simple, and comparatively few quantities have to be measured. There is, however, a sufficient number of variables to permit measuring inductances of a very wide range of values with the same bridge, using comparatively few values of the capacity.

(d) The method is particularly well adapted to measure inductance by the substitution method, where the inductance to be determined is replaced by a standard of nearly equal value. The difference between them can then be measured with very great precision, the residual errors of the bridge being nearly if not entirely eliminated.

There are no disadvantages of the method that are not shared by other methods, except so far as the use of a condenser may be deemed a disadvantage. There are, however, some sources of error to be guarded against which we shall discuss later.

3. ADVANTAGES AND DISADVANTAGES OF A VIBRATION GALVANOMETER.

When the bridge is completely balanced (the conditions for a resistance balance and an inductance balance being simultaneously satisfied) the current will be zero in the galvanometer at every instant. If, however, the steady current balance is slightly disturbed by the heating of the resistances, especially that of the inductance coil to be measured, no adjustment of the variable resistance r will make the current in the galvanometer zero. The result is that the needle of the galvanometer will have a certain minimum amplitude of vibration when r is correctly set. If now one of the resistances (say Q) is slightly altered, a complete balance may be attained and the needle will be perfectly still. This will be seen to be a distinct advantage, for one is always certain, when the needle is quiet, that *both of the conditions of the bridge are satisfied;* namely, the condition of the simple Wheatstone bridge ($P\,S = R\,Q$), and the condition imposed by the presence of the inductance which requires a particular value for the resistance r. But the vibration galvanometer does more than merely save the trouble of going back to the use of a direct current and a direct current galvanometer to see whether the balance still holds; for, when an appreciable current is used, the resistance may be changing sufficiently to render such a test insufficient. The vibration galvanometer, on the other hand, insures that at the very moment when the inductive balance is attained the resistance balance also holds, and thus no error from this cause can enter.

Still further, if the resistance of the inductive coil, or of the arms of the bridge, is different when carrying alternating current from its resistance when carrying direct current (as it always is, although the difference is very small for fine wires and low frequencies), the vibra-

tion galvanometer takes account of the true resistance under the con-
ditions of the experiment as a direct current galvanometer could not
do. This is of considerable importance in measuring the inductance
of coils of large wire. Neither a telephone nor an alternating current
d'Arsonval galvanometer possesses this advantage.

In practice it is not necessary to make a close adjustment of the
direct-current balance at all, as this can be determined just as well by
the vibration galvanometer. In our work a graduated scale is viewed
in a telescope by reflection from the mirror of the vibration galvano-
meter, the filament of an incandescent lamp used to illuminate the
scale being also seen in the telescope. When an approximate adjust-
ment of r and Q is secured, the filament will appear somewhat broad-
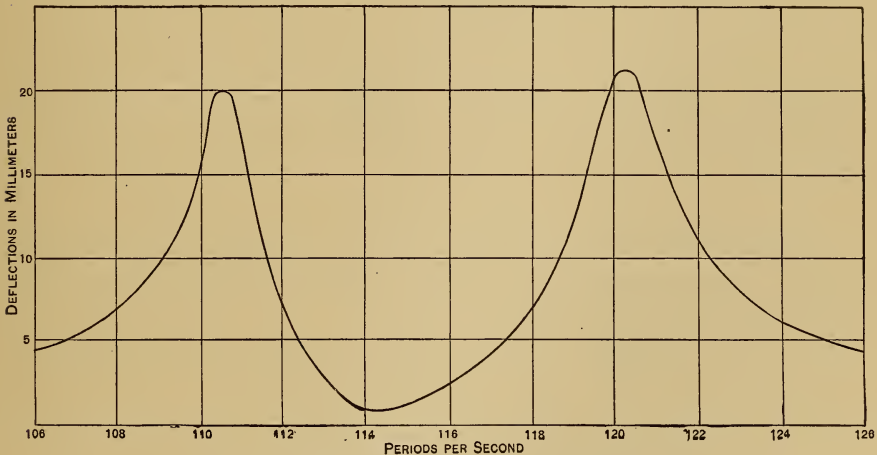
FIG. 5.—Sensibility curve of the vibration galvanometer.

ened by the slight vibration of the needle of the galvanometer. Small
changes in r and Q are then made successively until the filament appears
as a fine line and the lines on the scale are perfectly distinct[a]. This
adjustment can be made so delicately that a change in r or Q of one
part in a hundred thousand can be detected, when measuring induc-
tances of large values.

The chief disadvantage of the vibration galvanometer lies in the fact
that its sensibility decreases rapidly when the frequency of the cur-
rent varies from the natural period of the galvanometer. The sensi-
bility is nearly constant for a range of about one-half per cent in the
frequency but falls off rapidly when the frequency goes beyond this
range.

[a] Wien: Ann. d. Phys., **309**, p. 441; 1901.

In order to maintain the frequency at the point of maximum sensibility a Maxwell bridge is employed, as when measuring the capacity of a condenser. The condenser capacity and resistances of the bridge remaining constant, any change in speed causes a deflection of the galvanometer. An adjustable carbon resistance in the armature circuit of the driving motor permits the speed to be adjusted so that the deflection is reduced to zero. The motor is driven by current from a storage battery, and hence the changes in speed are relatively small. A glance at the galvanometer scale at any time shows whether the speed is correct, and if not, it is quickly adjusted by means of the rheostat.

Fig. 5 gives the sensibility curve of the vibration galvanometer, showing two peaks of high sensibility at 110.6 and 120 vibrations per second, respectively. At a frequency of 115 the sensibility is very low—much less than it is at frequencies outside the peaks of maximum sensibility. The curve is affected by changes of temperature, and can be altered at pleasure by varying the length and tension of the suspension wire.

4. THE APPARATUS.

A Rubens vibration galvanometer,[a] having a resistance of 200 ohms, is used. Its frequency may be varied between 100 and 200 per second, but has been used chiefly at about 110.

The several resistances are of manganin, and are all submerged in oil, to prevent heating and to enable their temperatures to be more accurately determined. The values of these resistances have been carefully measured every day that measurements of inductance have been made, when results of the highest accuracy have been sought. In series with the resistances r and Q, and forming part of them, are two slide wires which enable these resistances to be adjusted to 0.001 ohm, or even less, when necessary.

In order to eliminate as far as possible the errors due to slight changes in the arms P and R of the bridge, as well as any difference in their residual inductance and capacity, these resistances are always made equal and a commutator is employed to reverse them; a pair of readings is taken in every case, the mean of which is used in the calculation. The resistances Q and S were taken from two resistance boxes, in which the higher coils are subdivided to reduce the electrostatic capacity of the coil. We found in some of our early work that the residual capacity or inductance of noninductive resistances may be considerable; in the lower resistance coils the inductance predominates, and in the higher coils the capacity predominates. The connecting

---

[a] W. Oehmke, maker, Berlin.

Table I.—Results of Measurements of Two Inductance Coils of 1 Henry each; Coils Measured Singly and in Series, March 23, 1905.

[$P = R = 250.002$ ohms.]

1	2	3	4	5	6	7	8	9	10	13	14
No.	Coils.	S	r position 1.	r position 2.	r mean corrected.	$2r + P$	Temp. of condenser.	C	L measured values.	L_s sums of single values.	$\Delta = L_s - L$
		Ohms.	Ohms.	Ohms.	Ohms.	Ohms.	°	Microfarads.	Henrys.	Henrys.	Henrys.
1	F	500.204	868.62	868.45	868.53	1987.06	.....	1.00510	0.99900	.....	.....
2	C	"	883.72	883.56	868.64	2017.28	19.25	1.00510	1.01420	.....	.....
3	F	600.268	702.92	702.80	702.86	1655.72	.....	1.00511	0.99896		.....
4	C	"	715.565	715.435	715.50	1681.00	19.2	1.00511	1.01421		.....
5	F+C	"	1543.81	1543.56	1543.62	3337.24	.....	1.00511	2.01345	2.01317	−0.00028

Table II.—Results of Measurements of Four Inductance Coils of 100 Millihenrys each; Coils Measured Singly and in Series in Groups of Two, March 23, 1905.

[$P = R = 250.002$ ohms.]

1	2	3	4	5	6	7	8	9	10	13	14
No.	Coils.	S	r position 1.	r position 2.	r mean corrected.	$2r+P$.	Temp. of condenser.	C	L measured values.	L_s=sums of single values.	$\Delta = L_s - L$
		Ohms.	Ohms.	Ohms.	Ohms.	Ohms.	°	Microfarads.	Millihenrys.	Millihenrys.	Millihenrys.
1	A	250.102	74.000	73.972	73.999	398.000	19.0	1.00515	100.052		
2	B	"	73.904	73.872	73.901	397.802	18.95	1.00515	100.003		
3	H	"	73.861	73.822	73.855	397.712	19.0	1.00514	99.980		
4	C	"	73.880	73.846	73.876	407.754	19.0	1.00513	102.503		
5	H + C	250.102	277.737	277.665	277.708	805.418	19.1	1.00512	202.468	202.483	+ 0.015
6	A + B	"	272.907	272.832	272.872	795.746	19.1	1.00512	200.036	200.055	+ 0.019
										Mean..	+ 0.017
7	A + B	200.068	372.526	372.417	372.474	994.950		1.00511	200.075	200.092	+ 0.017
8	H + C	"	378.535	378.434	378.487	1006.976		1.00511	202.495	202.506	+ 0.011
										Mean..	+ 0.014
9	A	200.068	123.824	123.775	123.812	497.626		1.00510	100.067		
10	B	"	123.718	123.673	123.708	497.418		1.00510	100.025		
11	H	"	123.664	123.614	123.650	497.302		1.00510	100.002		
12	C	"	129.886	129.836	129.872	509.746	19.25	1.00510	102.504		

leads have substantial terminals and the resistances of these leads, in thousandths of an ohm, is stamped on the terminals.   Their values are always included in making up the corrected resistances of the bridge.

The measurements shown in Tables I and II, made March 23, illustrate the results obtained in the determination of inductances of one henry and one-tenth henry.   An electromotive force of about 50 volts was employed on the bridge.

Fig. 6 shows how the commutator was connected to the bridge so as to reverse P and R, which are equal.   Formula (2) in this case ($Q = S$) reduces to

$$L = C\,S\,(2r + P).$$

Columns 4 and 5 give the nominal values of r in the two positions of the commutator, and column 5 the mean value, corrected from the
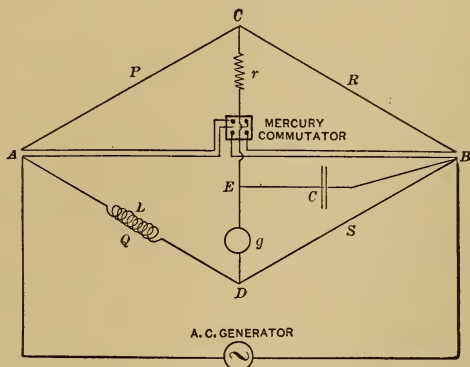
FIG. 6.—Showing commutator for interchanging two
arms of the Anderson Bridge.

results of the latest comparison of the resistances with standard resistances.

Column 9 gives the capacity of the condenser C.   Where two inductances are measured in series the measured sum is given in column 10, and the sum of the separate values are given in the next column.   The last column gives the differences between these measured values and the sums of the separate values.   While these differences are very small, averaging less than one in ten thousand, they are appreciable and always positive.   This indicates that there may be some constant source of error in the bridge.

5. SOURCES OF ERROR.

The results given above show that measurements of inductance of very great precision can be made by Anderson's method, provided there are no constant errors of appreciable magnitude entering into

the results.   Such errors might be due to any one of the following causes:

(a) The residual inductance or capacity of the resistance r and of the arms of the bridge (not including, of course, the inductance of the coil in Q which is being measured) may introduce a constant error in L.   As stated above, we have always made the arms P and R equal in value and reversed them by a commutator, in order to eliminate any difference there may be in their resistances or in their inductances or capacities.   But the differences between Q and S can not thus be eliminated.   The total resistance of Q is, of course, equal to S, since $P=R$ (except for a small change due to residual inductances, to be discussed later), but part of this is in the inductive coil itself.   Residual inductance in the noninductive part of Q makes the measured value of L too large, and, conversely, capacity would make it too small.   The effect of inductance or capacity in S is of opposite sign to that of Q, and hence if the resistances of Q and S are *similar*, that is, made up as far as possible of the same kind of coils, then their effects will balance except for that part of S which equals the resistance of the inductive coil.   In our work S was fixed in any given case, and Q was varied to secure a balance; thus Q usually contains a number of small resistances in addition to the slide wire, and these can not be counterbalanced exactly by S.

(b) The inductive coils must be removed some distance from the bridge and from each other when two or more are measured at once. This requires leads of one to three meters in length (for the larger inductances), and these leads may affect the measured value of the inductance.   If they are close together, so as to be noninductive (as twisted lamp cord, for instance), they possess an appreciable capacity; and if far enough apart to be free from capacity, they possess measurable inductance.   In measuring small inductance coils the capacity effect is small, and it is better to have the leads close together and as short as is safe.   With large inductances the capacity of the leads is more important, and it is better to have them farther apart, to reduce it to a minimum.   The inductance of the leads can then be calculated (or separately measured) and applied as a correction, if desired, or the same leads may always be employed with a given coil and considered as a part of the coil.   The inductance of the wires joining the condenser to the bridge tends to reduce the capacity in the ratio of pl to $\dfrac{1}{pc}$ or p^2lc to unity, where l is the small inductance of the leads; on the other hand, if these leads are close together, their capacity is added to that of the condenser.   In our experiments, where the leads

were short and wide apart, both these effects were inappreciable.   But
if currents of high frequency are used, particularly with large capacity,
the error due to the inductance of the leads may be appreciable; on
the other hand, with small condenser capacity, the error due to the
capacity of lead wires near together (as a twisted lamp cord) may be
considerable and of opposite sign to the other.   In precision measure-
ments, therefore, care should be taken that no error is introduced in
this manner.

(c) The inductive coil itself has a certain electrostatic capacity which
modifies its measured inductance by an amount depending on the fre-
quency of the current and the inductance of the coil, as well as its
capacity.[a]   The approximate value of the expression for the measured
inductance L' in terms of the true inductance L is

$$L' = L\,(1 + p^2 cL) \tag{4}$$

where c is the electrostatic capacity of the coil and $p = 2\pi$ times the
frequency.   In practice, c is found by measuring L' at two different
frequencies; it is too small to be important for the smaller values of
inductance at low frequencies.   One of our inductance coils, having
an inductance of 1 henry, has a capacity of 1×10^{-10} farads.   For a
frequency of 112, this value makes the correction term $p^2 CL$ in the
above expression .00005, a quantity which can be detected, but which
is not a large error.   If, however, the frequency were ten times as
great, this term would become .005, a very important correction.

The electrostatic capacity of a coil can be made relatively small by
winding it in a deep channel, so that there are many layers and com-
paratively few turns in a layer.   This, however, reduces its induct-
ance, and in practice it is better not to depart very far from the form
giving maximum inductance.   The electrostatic capacity of the cord,
as already pointed out, increases the value of this correction term.

(d) The capacity of the condenser, as already stated, can be deter-
mined with very great accuracy, and by taking careful account of the
temperature of the condenser and its temperature coefficient, there
will be very little uncertainty in the value of the capacity.   The ques-
tion remains, however, as to what effect the absorption in the con-
denser produces on the measured value of the inductance when used
in Anderson's method.   The effect of absorption is to cause the
current to lag a little behind its phase in a perfect condenser.   That
is, it is in advance of the electromotive force by a little less than 90°.
We give below a theoretical investigation of this question, and also

[a] Wien: Ann. d. Phys., 44, p. 711; 1891.   Dolezalek: Ann. d. Phys., 12, p. 1153;
1903.

experimental measures, wherein the phase of the condenser current is shifted back by placing a resistance in series with it.   Fortunately, the error due to the slight displacement of phase produced by the small absorption in a good mica condenser, is inappreciable.   The effect of slight leakage is also investigated below, and proves to be inappreciable.

6. CALCULATION OF THE EFFECT OF THE RESIDUAL INDUCTANCES AND CAPACITIES OF THE ARMS OF THE BRIDGE.

Inductance in any arm of the bridge causes the current to lag, while capacity advances its phase.   The angular lag due to an inductance l is $\dfrac{pl}{R}$, and the angular advance due to capacity is pcR.   The latter value follows from the fact that the capacity c is in parallel with the resistance.   The current through the resistance is $\dfrac{E}{R}$; the capacity current $90°$ ahead of this is pcE.   The ratio of these two currents is pcR, and this is the angle of advance of the resultant current.   If $\dfrac{pl}{R} = pcR$, the current will have the same phase as though both capacity and inductance were absent.   Thus, if $l = cR^2$ the coil may be considered free from inductance; if $l < cR^2$, it may be considered to possess negative inductance equal to $cR^2 - l$.
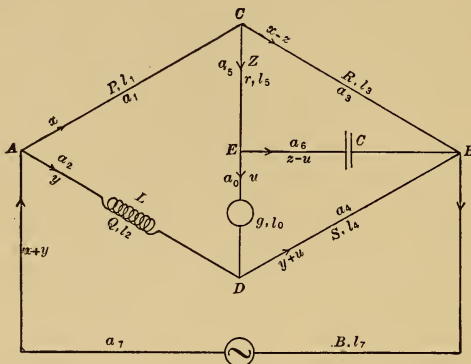
Fig. 7.—Network of Anderson Bridge, with currents, resistances, inductances, and impedances of the arms separately indicated.

The actual values of these residual positive or negative inductances vary widely, according to the length of wire on a coil, and its size, thickness of insulation, and manner of winding.   High-resistance coils, made of many turns of fine wire wound in the usual noninductive manner, have relatively large capacity, which we here regard as negative inductance.   In low-resistance coils of fewer turns of coarse wire the inductance predominates, and l is therefore positive.

To calculate the effect of these residual inductances, we shall form the equations for the networks of the Anderson bridge, assuming each branch to have positive or negative inductance, and solve for L the inductance of the coil to be measured. In fig. 7 the resistances, inductances, impedances, and currents in each arm of the bridge are indicated, and the positive directions of the currents are shown by arrows. Thus,

$g, P, Q, R, S, r, 0, B$ are the resistances.
$l_0, l_1, l_2, l_3, l_4, l_5, 0, l_7$ are the small inductances, $+$ or $-$.
$a_0, a_1, a_2, a_3, a_4, a_5, a_6, a_7$ are the impedances.
$u, x, y, x - z, y + u, z, z - u, x + y$ are the currents.

The values of the impedances in the eight branches of the bridge are then as follows:

$$\left.\begin{aligned}
a_0 &= g + i\,p\,l_0 \\
a_1 &= P + i\,p\,l_1 \\
a_2 &= Q + i\,p\,(l_2 + L) \\
a_3 &= R + i\,p\,l_3 \\
a_4 &= S + i\,p\,l_4 \\
a_5 &= r + i\,p\,l_5 \\
a_6 &= 0 + \frac{1}{ip\,C} \\
a_7 &= B + i\,p\,l_7
\end{aligned}\right\} \qquad (5)$$

Applying Kirchhoff's law to the four circuits $A\,C\,D$, $C\,B\,E$, $E\,B\,D$, $A\,D\,B$, we have the following equations:

$$\left.\begin{aligned}
a_1 x + a_5 z + a_0 u - a_2 y &= 0 \\
a_3 (x - z) - a_6 (z - u) - a_5 z &= 0 \\
a_6 (z - u) - a_4 (y + u) - a_0 u &= 0 \\
a_7 (x + y) + a_2 y + a_4 (y + u) &= E
\end{aligned}\right\} \qquad (6)$$

Rearranging the terms of these equations, we have

$$\left.\begin{aligned}
a_0 u + a_1 x - a_2 y + a_5 z &= 0 \\
a_6 u + a_3 x - (a_3 + a_5 + a_6)z &= 0 \\
(a_0 + a_4 + a_6)u + a_4 y - a_6 z &= 0 \\
a_4 u + a_7 x + (a_2 + a_4 + a_7)y &= E
\end{aligned}\right\} \qquad (7)$$

Solving for u, the current through the galvanometer, we obtain

$$u = \frac{1}{\varDelta}\begin{vmatrix} 0 & a_1 & -a_2 & a_5 \\ 0 & a_3 & 0 & -(a_3 + a_5 + a_6) \\ 0 & 0 & a_4 & -a_6 \\ E & a_7 & (a_2 + a_4 + a_7) & 0 \end{vmatrix} = -\frac{E}{\varDelta}\begin{vmatrix} a_1 & -a_2 & a_5 \\ a_3 & 0 & -(a_3 + a_5 + a_6) \\ 0 & a_4 & -a_6 \end{vmatrix} \qquad (8)$$

If $u=0$, we then have

$$a_1\, a_4\, (a_3+a_5+a_6)-a_3\, a_2\, a_6+a_3\, a_4\, a_5=0$$

or,

$$a_1\, a_3\, a_4+a_1\, a_4\, a_5+a_1\, a_4\, a_6+a_3\, a_4\, a_5-a_2\, a_3\, a_6=0 \qquad (9)$$

Substituting the values of the impedances given in (5) above, we have

$$\left.\begin{aligned}
&(P+ipl_1)\,(R+ipl_3)\,(S+ipl_4)+(P+ipl_1)\,(S+ipl_4)\,(r+ipl_5)\\
&+(P+ipl_1)\,(S+ipl_4)\frac{1}{ip\,C}+(R+ipl_3)\,(S+ipl_4)\,(r+ipl_5)\\
&-(Q+ip(l_2+L))\,(R+ipl_3)\frac{1}{ip\,C}=0
\end{aligned}\right\} \quad (10)$$

Separating the real and imaginary parts, we have first, for the real part

$$PRS+PrS+RSr-\frac{LR}{C}+\frac{1}{C}\Big[\,l_1S-l_2R-l_3Q+l_4P\,\Big]$$
$$-p^2\Big[\,rl_4(l_1+l_3)+S(l_1l_3+l_1l_5+l_3l_5)+Rl_4(l_1+l_5)+Pl_4(l_3+l_5)\,\Big]=0$$

Transposing and dividing by $\dfrac{R}{C}$,

$$\left.\begin{aligned}
L=CS\Big[\,r\Big(\frac{P+R}{R}\Big)+P\,\Big]+\frac{1}{R}[l_1S-l_2\,R-l_3Q+l_4\,P]-\\
\frac{p^2\,C}{R}\Big[\,rl_4\,(l_1+l_3)\\
+S\,(l_1\,l_3+l_3\,l_5+l_5\,l_1)+Rl_4\,(l_1+l_5)+Pl_4\,(l_3+l_5)\Big]
\end{aligned}\right\} \quad (1^1)$$

or, if $L_0=CS\Big[\,r\Big(\dfrac{P+R}{R}\Big)+P\,\Big]$,

$$L=L_o+\alpha-\beta \qquad (12)$$

If the small inductances l_1, l_2, l_3, l_4, l_5 are all zero, as in the ideal case, the last two terms disappear and

$$L=C\,S\Big[\,r\frac{(P+R)}{R}+P\,\Big] \qquad (13)$$

and since in the Wheatstone bridge $\dfrac{P}{R}=\dfrac{Q}{S}$ we have

$$L=C\,[r(Q+S)+PS],$$

which is the expression (2) given above for the inductance by Anderson's method.

The last term β in equation (12), having a coefficient $\dfrac{p^2\,C}{R}$ is negligible, unless the frequency is very high or the residual inductances excessive.

The second term α gives the principal correction due to the induct-ances in the four arms of the bridge. It will be seen that the induct-ance of the resistance r enters only in β, which is negligible when l_5 is small, and the galvanometer inductance has disappeared entirely from the equation. The second term α consists of four parts, two of which are positive and two negative. The part $\frac{1}{R}(l_1 S - l_3 Q)$, due to the two arms P and R, is eliminated in our experiments by making $P = R$ and reversing P and R. The remainder $\frac{1}{R}(l_4 P - l_2 R) = l_4 - l_2$ (since $P = R$) is due to the two remaining arms Q and S. As stated above, if these two arms consist of resistances made up of similar coils, (i. e., coils of the same resistance wound in the same manner with the same size wire) their inductances will be nearly equal. They might be exactly equal, if the inductive coil L had no resistance. It is therefore desirable to have l_2 and l_4 as nearly equal as possible, and then when necessary apply a correction for their difference.

If we divide the expression for α in equation (11) above, by Q, remembering that $Q = \frac{PS}{R}$ (approximately), we obtain

$$\left. \begin{aligned} \alpha &= Q\left[\frac{l_1}{P} - \frac{l_2}{Q} - \frac{l_3}{R} - \frac{l_4}{S}\right] \\ &= \frac{Q}{p}\left[\frac{pl_1}{P} - \frac{pl_2}{Q} - \frac{pl_3}{R} + \frac{pl_4}{S}\right] \\ &= \frac{Q}{p}\left[\phi_1 - \phi_2 - \phi_3 + \phi_4\right] \end{aligned} \right\} \quad (14)$$

where ϕ_1, ϕ_2, ϕ_3, ϕ_4 are the phase angles of the currents in the four arms of the bridge, due to the combined inductance and capacity of the resistances P, Q, R, S, neglecting, of course, the inductance L in Q, which is to be measured. These angles may be positive or negative. If they are all equal, or if $\phi_1 + \phi_4 = \phi_2 + \phi_3$ (algebraically) the correction term reduces to zero.

As we shall show below, these angles are appreciable in the "nonin-ductive" windings usual in resistance boxes, and the correction α is therefore important in precision work. The resistances may, how-ever, be so wound and adjusted as to make the angles ϕ inappreciable.

The imaginary part of equation (10) above gives

$$\left. \begin{aligned} PS - RQ &= p^2(l_1 l_4 - l_3(l_2 + L)) + p^2 C[PRl_4 + SR(l_1 + l_5) \\ &+ PS(l_3 + l_5) + r(Sl_1 + Sl_3 + Pl_4 + Rl_4)] - p^4 C(l_1 l_3 + l_3 l_5 + l_5 l_1) l_4 = \gamma \end{aligned} \right\} (15)$$

If l_1, l_2, l_3, l_4, l_5 are all zero, this reduces to $PS - RQ = 0$, which is the

condition for the Wheatstone bridge and the condition assumed in the ideal Anderson bridge. But when these quantities are not zero this condition does not hold and Q is not equal to $\dfrac{PS}{R}$. Consequently, an error is introduced in calculating L from the formula (2) of Anderson, unless Q is assumed equal to $\dfrac{PS}{R}$, instead of using its actual value.

The variation of Q, due to changes in one or more of the small inductances l (P, R, and S remaining unchanged), is illustrated in some of the examples given below.

7. ILLUSTRATION OF THE FORMULÆ.

In order to ascertain the order of magnitude of the corrections α and β of formula (12), and the variation of $PS-RQ$ from zero in (15), we shall assume the following values of the constants for cases I, II, and III:

$$L = 1 \text{ henry,}$$
$$C = 1 \text{ microfarad,}$$
$$P = R = 250 \text{ ohms,}$$
$$S = 500 = Q \text{ (approximately),}$$
$$r = 875,$$
$$p^2 = 500{,}000.$$

INDUCTANCES IN MICROHENRYS.

Case.	I	II	III	IV
l_1	− 2	+25	− 50	−1,500
l_2	+ 2	+40	+100	− 250
l_3	+ 2	+25	+ 50	0
l_4	− 2	+50	−100	−5,000
l_5	+10	+50	+100	+ 100

Case I is supposed to represent a specially wound bridge in which all the inductances are small, but in order to make it represent the most unfavorable case two are taken positive and two negative, so as to give a maximum value to the error α.

Case II represents a favorable arrangement where the coils are similar and the inductances all positive, but with larger values than those of case I, being such as might be expected in practice.

In Case III we have assumed that P is a single coil of fine wire, having the capacity effect greater than the inductance by an amount

equivalent to a negative inductance of 50 microhenrys, whereas R is made up of several coils of coarser wire, giving a smaller capacity and larger inductance, and hence l_3 is taken as $+50$. Similarly, S is supposed to be a single coil of 500 ohms, with capacity predominating, and equivalent to a negative inductance of 100 microhenrys, while Q is the sum of several smaller coils, and l_2 is therefore positive and equal to 100 microhenrys. This is perhaps an extreme case, but not an improbable one.

In case IV larger resistances are assumed, viz:

$$P=1,000 = Q \text{ (approximately)}.$$
$$R=S=5,000.$$
$$r=833.3.$$
$$C=0.1 \text{ microfarad}.$$

P, Q, and S are assumed to have negative inductances, each being supposed to consist of a single coil (or mainly of large coils, as in the case of Q), while R is supposed to be made up of smaller coils having the inductance and capacity balanced. The values of the inductances chosen for P, Q, and S are approximately those found in noninductive resistances of such magnitudes.

Substituting the above values of the constants in equation (11) we find the following values of α and β:

Case.	L_0	α	β
I	1 Henry	-0.000012	$+ \ 4 \times 10^{-12}$
II	1 "	$+0.000010$	$- \ 1.0 \times 10^{-8}$
III	1 "	-0.000400	$+ \ 1.2 \times 10^{-8}$
IV	1 "	-0.002250	$-34 \ \times 10^{-8}$

These results show that the β term is small in comparison with α, and may be neglected. The correction α is as large in case I as in Case II, showing that if the several small inductances are all of the same sign, and proportional to the resistances, they cancel out, except for the necessary inequality in l_2 and l_4, due to the resistance of the coil to be measured. If the inductances of the coils are adjusted to as small values as 2 microhenrys for each arm, they may be either positive or negative without making the error appreciable, as the error in case I, with l_1 and l_4 opposite in sign to l_2 and l_3 is the greatest possible for such values of l_1, l_2, l_3, l_4.

The results shown in Table III illustrate the importance of the correction term α for coils of smaller inductances, namely: 100, 10, 1, 0.1, and 0.01 millihenrys.

Table III.—Showing the Values of the Correction α for Various Values of Inductances and the Corresponding Values of Capacities and Resistances.

Inductance to be measured.	Capacity.	$P = R$	l_1	l_3	$Q = S$	l_2	l_4	r	α
Millihenrys.	Microfarads.	Ohms.	Microhenrys.	Microhenrys.	Ohms.	Microhenrys.	Microhenrys.	Ohms.	Millihenrys.
100	1.0	250	+2	−2	250	−2	+2	75	0.008
10	0.4	100	+1	−1.0	100	−1	+1	75	0.004
1	0.1	50	+0.5	−0.5	100	−1	+1	25	0.004
0.1	0.05	20	+0.5	−0.5	50	−0.5	+0.5	10	0.0045
0.01	0.02	20	+0.5	−0.5	20	−0.5	+0.5	2.5	0.002

It will be seen that the error α in the case of 100 millihenrys is scarcely appreciable, but that in the others it is appreciable and in the smaller coils it amounts to several per cent. These computed errors, as before, are the maximum values for the assumed residual inductances, since we have taken l_1 and l_4 of opposite sign to l_2 and l_3. In practice we should therefore expect smaller errors on the average unless the values of the l's are larger. If the coils are not wound to a minimum value of the inductance, the errors may be much larger than those above. It is therefore our practice in measuring small inductances to take them by difference, leaving P, R, and S unchanged and altering Q to compensate for the resistance of the coil to be measured. If the inductances of the resistances of Q, or at least of the part to be replaced by the coil to be measured, are accurately known, the difference of two determinations gives the true inductance desired.

The value of $PS - RQ$ is found from equation (15) by substituting the values of the resistances and residual inductances. In case I above, $S - Q$ is only $-.003$ ohm, while in Case IV it is -4.55 ohms. In other words, Q is larger by 4.55 ohms in a total of 1,000 when the bridge is exactly balanced for alternating current than it is for a direct-current balance. Consequently, if, as is sometimes done, the bridge is balanced with direct current, and then an alternating current is applied, the resistance balance no longer holds, if there are residual inductances and capacities, and Q may require a change of several ohms to secure the resistance balance, the inductive balance being

effected as we have seen above by varying r. In calculating L, how-
ever, no account need be taken of Q, as it is eliminated from the
expression (13) for L. Hence there is no occasion to calculate the
value of γ in (15).

8. EFFECT OF RESISTANCE IN SERIES AND IN PARALLEL WITH THE CONDENSER.

As a mica condenser is not entirely free from absorption and might
also show slight leakage, it is desirable to ascertain how large an error,
if any, is produced by using such a condenser instead of the ideal con-
denser assumed in the theory of the Anderson bridge, namely, one in
which the impedance is $\dfrac{1}{ipC}$. Resistance in series with a perfect con-
denser produces the same phase displacement as a certain amount of
absorption, and resistance in parallel with the condenser has the effect
of leakage or imperfect insulation. In a good mica condenser the
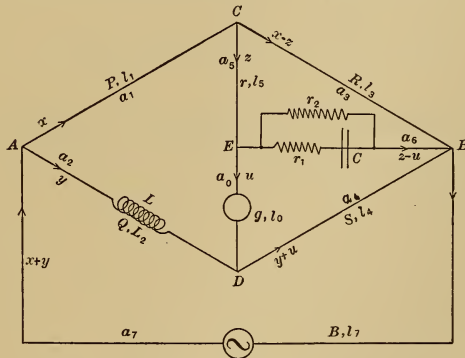
FIG. 8.—Resistance in series and in parallel with the condenser.

phase of the current with a frequency of 100 per second should dif-
fer from quadrature with the electromotive force by not more than
one minute of angle, and may be as small as 30″, although it may be
several minutes in inferior condensers (even as high as 30′). For paper
condensers the angle may be as small as 4′ and as large as several
degrees. For a condenser of one microfarad capacity, with a frequency
of 100 per second, 30″ of angle corresponds to a resistance in series
with the condenser of 0.23 ohm, whereas 30′ would correspond to a
resistance of 14 ohms.

In other words, such resistances in series with perfect condensers
would give currents of the same phases as the imperfect condensers
employed. A leakage resistance less than a thousand megohms would
never occur in a good condenser. We shall now calculate (1) the effect

of introducing resistance in series with the condenser to correspond with absorption and (2) of placing a high-resistance shunt around the condenser to represent leakage. In fig. 8 these two resistances are represented by r_1 and r_2. If we substitute the values of the impedances of the arms of the bridge in formula (9), we shall obtain an expression for the inductance to be measured in the same manner as before. In this case, however, we assume for convenience that the resistances P, Q, R, S, and r are all free from inductance or capacity (except, of course, L in Q), and hence

$$a_1 = P$$
$$a_2 = Q + ipL$$
$$a_3 = R$$
$$a_4 = S$$
$$a_5 = r$$
$$a_6 = r_1 + \frac{1}{ipC} \text{ in the first case,}$$

and
$$a_6 = \frac{r_2}{1 + ipCr_2} \text{ in the second case.}$$

(a) *Resistance in series with the condenser.*—Equation (9) becomes, when the above values of the impedances are substituted,

$$PRS + PSr + PS\left(r_1 + \frac{1}{ipC}\right) + RSr - (Q + ipL)\left(r_1 + \frac{1}{ipC}\right)R = 0 \quad (16)$$

Separating the real and imaginary parts, we have, first, for the real part,

$$PS(R + r + r_1) + RSr - QRr_1 - \frac{L}{C}R = 0$$

or,
$$L = CS\left[r\frac{(P+R)}{R} + P\right] + Cr_1\left[\frac{PS}{R} - Q\right] \quad (17)$$

The first term of this expression is the same as that of (13), and is the value of L when r_1 is zero. $\frac{PS}{R} - Q$ would be zero in the ideal bridge. To find its value in this case we make use of the imaginary part of (16) above.

This imaginary part is

$$\frac{PS - RQ}{ipC} - ipLRr_1 = 0$$

or,
$$PS - RQ = -p^2LRr_1C.$$

Whence
$$Q = \frac{PS}{R} + p^2Lr_1C. \quad (18)$$

Substituting the value of $\dfrac{PS}{R} - Q$, derived from (18), in equation (17), we have

$$L = CS[r\frac{(P+R)}{R} + P] - p^2 r_1^2 L C^2 \qquad (19)$$

In fig. 9 ,which is the impedance diagram of a condenser, r_1 is the series resistance, equivalent (in its effect upon the phase angle) to the absorption, and θ is the small angle by which the current falls short of 90° in its phase relation to the impressed electromotive force. Hence tan $\theta = p C r_1$. Substituting in (19) above,

$$L = L_0 - L \tan^2 \theta$$

or, $$L = L_0 (1 - \tan^2 \theta) \qquad (20)$$

FIG. 9.—Impedance diagram of an imperfect condenser.

since L_0, the value of L when the correction is zero, is substantially the same as L. In the best mica condensers, as stated above, θ is about half a minute, and tan θ is 0.00015; $\tan^2 \theta$ is therefore only about two parts in a hundred million. Hence the angle θ might be ten times as large without producing an appreciable error, although in some mica condensers that we have tested the angle is large enough to produce a sensible error.

(b) *Resistance in parallel with the condenser.*—Substituting $\dfrac{r_2}{1 + ip C r_2}$ for a_6 in equation (9) above, we get a solution for the case of resistance in parallel with the condenser. The direct substitution gives

$$PRS + PSr + PS\left(\frac{r_2}{1 + ip C r_2}\right) + RSr - R(Q + ipL)\left(\frac{r_2}{1 + ip C r_2}\right) = 0 \,(21)$$

The real part of this is

$$PRS + PSr + RSr + (PS - RQ)r_2 = 0$$

Hence, $$Q - \frac{PS}{R} = \frac{S}{r_2}[r\frac{(P+R)}{R} + P]$$

or, $$Q - \frac{PS}{R} = \frac{L_0}{r_2 C} \qquad (22)$$

Thus, the variation in Q is inversely proportional to r_2, the shunt resistance, and to the capacity of the condenser, and directly proportional to L. The leakage resistance through a condenser is inversely proportional to the capacity, so that in general $r_2 C$ is independent of the value of the capacity, but depends on the quality of the condenser.

If r_2=1000 megohms and C=1 microfarad, $r_2 C$=1000 and the variation of Q is in that case very small.

The imaginary part of equation (21) above is

$$Cr_2(PRS+PSr+RSr)=Lr_2R$$

or, $$L= CS\,[r\frac{(P+R)}{R}+P]=L_\circ \qquad (23)$$

This is equation (13), and shows that r_2 *has no effect whatever on the measured value of* L.

9. VERIFICATION OF FORMULÆ 11, 18, 19, 22, 23.

In Table IV the results are given of a series of measurements made to verify formula (11) by introducing a small inductance successively in each of the arms of the bridge; the β term in this formula is negligible, in comparison with α, except for the case of inductance in r only, in which case α is zero and the β term has then been computed. This is the fifth case in the table, where the observed change ($\Delta\ \beta$ in this case, as $\Delta\alpha$=o) is 0.01 millihenry, whereas the calculated effect is still smaller. But the inductance coil measured has an inductance of 1 henry, and hence the observed change is only one part in a hundred thousand, a quantity barely measureable.

If we differentiate formula (12) we have, since L does not change when the inductance coil is inserted in any arm of the bridge, and β is assumed zero,

$$\Delta L_\circ+\Delta\alpha=o$$

or, $$\Delta\alpha= -\Delta L_\circ$$

where $\Delta\alpha$ is the change produced in the α term by the insertion of the given coil in any arm of the bridge, and $\Delta L_\circ$ is the change in the first term of equation (11) or (12); that is, in the computed value of L apart from the correction terms.

For example, in the first case, 0.5 millihenry was inserted in the arm Q, and hence l_2 was increased by 0.5 millihenry. The second term of the α correction was therefore increased by $0.5\times\frac{R}{R}$=0.500 millihenry. The observed value of $\Delta L_\circ$ in this case was the same. In case 3, the addition of 0.5 millihenry in P changed the first term of the α correction by $0.5\times\frac{S}{R}$=1.000 millihenry, whereas the measured change was 0.999 millihenry. In the fourth case the difference

Table IV.—Effect of placing an Inductance Coil of 0.5 Millihenry in the Arms of the Anderson Bridge.

[$P.=R=250$ ohms. $S=500$ ohms. $C=1$ microfarad. $L=1$ henry.]

1	2	3	4	5	6	7	8	9	10	11
No.	Position of coil.	r position 1.	r position 2.	r mean.	Δr	ΔS	ΔR	ΔP	ΔL_o observed.	ΔL_o calculated.
		Ohms.	Ohms.	Ohms.	Ohms.	Ohms.	Ohms.	Ohms.	Milli-henrys.	Milli-henrys.
1	Coil in Q ...... Coil removed	884.500 884.001	884.330 883.836	884.415 883.918	+0.497	0	0	0	+0.500	+0.500
2	Coil in S ...... Coil removed	881.911 884.025	881.741 883.853	881.826 883.939	−2.113	+0.793	0	0	−0.519	−0.500
3	Coil in P ...... Coil removed	881.270 884.051	881.084 883.883	881.177 883.967	−2.790	0	0	+0.794	−0.999	−1.000
4	Coil in R ...... Coil removed	886.439 884.051	886.268 883.883	886.354 883.967	+2.387	0	+0.794	0	+0.989	+1.000
5	Coil in r ...... Coil removed	884.051 884.040	883.886 883.878	883.969 883.959	+0.010	0	0	0	+0.010	+0.0002

between the observed and calculated value was larger, namely, 0.011 millihenry. This is, however, only 11 parts in a million compared with the coil being measured, and is a very small discrepancy. These results may be regarded as fully verifying the formula when the inductances are positive.

Table V.—Effect of placing a Capacity of 0.00945 Microfarad in Parallel with the Arms of the Anderson Bridge.

[P=R=250 ohms. S=500 ohms. C=1 microfarad. L=1 henry.]

1	2	3	4	5	6	7	8
No.	Position of the condenser.	r position 1.	r position 2.	r mean.	$\varDelta\,r$	$\varDelta\,L_0$ observed.	$\varDelta\,L_0$ calculated.
		Ohms.	Ohms.	Ohms.	Ohms.	Milli-henrys.	Milli-henrys.
1	Condenser around Q.	882. 265	882. 065	882. 165	—1. 529	—1. 536	—1. 550
	Condenser removed .	883. 793	883. 596	883. 694			
2	Condenser around S .	886. 195	885. 930	886. 062	+2. 351	+2. 367	+2. 362
	Condenser removed .	883. 827	883. 595	883. 711			
3	Condenser around P.	884. 997	884. 793	884. 895	+1. 171	+1. 176	+1. 181
	Condenser removed .	883. 825	883. 623	883. 724			
4	Condenser around R	882. 618	882. 435	882. 527	—1. 183	—1. 188	—1. 181
	Condenser removed .	883. 820	883. 600	883. 710			
5	Condenser around r .	883. 860	883. 671	883. 766	+0. 047	+0. 047	—0. 0004
	Condenser removed .	883. 818	883. 620	883. 719			

In the columns headed $\varDelta S$, $\varDelta R$, $\varDelta P$, the quantity 0.794 represents the resistance of the inductance coil of 0.5 millihenry inserted in the arms.

To test the formula for the case of negative inductances, a capacity was inserted in parallel with each of the five arms of the Anderson bridge in succession, and the changes in r observed. From these changes $\varDelta L_0$ was determined and the result compared with the computed change in the α correction term. As stated above, a capacity C placed in parallel with a resistance R is equivalent to a negative inductance l, determined by the expression

$$l = CR^2$$

This capacity may be located in the resistance coil itself or in a condenser joined in parallel with it.

In the first case of Table V the resistance of Q was 405 ohms noninductive and 95 ohms in the coil whose inductance was 1 henry. The condenser was placed in parallel with the former, and had an effect proportional to $\overline{405}^2$ as compared with an effect proportional to $\overline{500}^2$ in S, given in case 2. It will be seen that the differences between the observed and calculated changes, due to capacity, is only a few thousandths of a millihenry—that is, only a few parts in a million of the coil each time measured. Hence the formula may be regarded as completely verified for negative inductances as well as positive.

Table VI.—Effect of Placing Resistance r_1 in Series With the Condenser of the Anderson Bridge.

[$P=R=250$ ohms. $S=250$ ohms. $L=1$ henry. $C=2$ microfarads. $p^2=474{,}800$. $n=110$ per second.]

1	2	3	4	5	6	7	8	9
r_1	r Position 1.	r Position 2.	r Mean.	$\varDelta L_0$	$p^2 r_1^2 C^2 L$	Q'	$\varDelta Q$	$p^2 r_1 CL$
Ohms.	Ohms.	Ohms.	Ohms.	Milli-henrys.	Milli-henrys.	Ohms.	Ohms.	Ohms.
0	871.82	871.59	871.705	--------	--------	155.27	--------	--------
5	871.89	871.68	871.785	+0.08	+0.05	--------	--------	--------
20	872.625	872.40	872.51	0.80	0.76	173.86	+19.43	+19.00
40	874.92	874.72	874.82	3.12	3.05	189.80	35.78	38.00
80	884.07	883.98	884.025	12.32	12.19	225.99	72.39	76.00
0	----------	----------	----------	--------	--------	153.18	--------	--------

Table VI gives the results of measurements made to verify formulæ (18) and (19). Resistances of 5, 20, 40, and 80 ohms were placed in

Table VII.—Effect of Placing Resistance r_2 in Parallel with the Condenser of the Anderson Bridge.

[$P=R=250$ ohms. $S=250$ ohms. $C=2$ microfarad. $L=1$ henry.]

1	2	3	4	5	6	7	8	9	10	11
r_2	r Position 1.	r Position 2.	r Mean.	ΔL_0	$\frac{L_0}{C}$	C calc.	Capacity in m.f., single coil (coil=100,000 ohms).	Q'	ΔQ	$\frac{L_0}{C}\cdot\frac{1}{r_2}$
Ohms.	Ohms.	Ohms.	Ohms.	Milli-henrys.	×10⁶	Microfarads.		Ohms.	Ohms.	Ohms.
∞	871.83	871.60	871.715		0.4984			153.18		
1,000,000	871.83	871.59	871.71	−0.005	0.4984	0.00012	0.00048	153.80	+0.62	+0.50
400,000	871.77	871.55	871.66	−0.060	0.4983	0.00026	0.00052	154.41	1.23	1.24
200,000	871.71	871.48	871.595	−0.130	0.4983	0.00059	0.00059	155.75	2.57	2.49
100,000	871.54	871.33	871.435	−0.295	0.4982	0.00124	0.00062	158.07	4.89	4.98
50,000	871.23	871.00	871.115	−0.615	0.4980	0.00251	0.00063	162.83	9.65	9.96
25,000	870.60	870.37	870.485	−1.250	0.4978	0.00517	0.00064	172.44	19.26	19.91
12,500	869.31	869.08	869.195	−2.550	0.4971			191.52	38.34	39.76
∞	871.86	871.63	871.745					153.18		
∞	871.87	871.665	871.77		0.4984		Coil=10,000 ohms.	152.39		
100,000	871.78	871.57	871.675	−0.095	0.4983	0.000191	0.00191	157.15	+4.76	+4.98
90,000	871.75	871.58	871.665	−0.105	0.4983	0.000211	0.00189	157.61	5.22	5.54
70,000	871.74	871.56	871.65	−0.125	0.4983	0.000251	0.00176	159.18	6.80	7.12
50,000	871.68	871.49	871.585	−0.190	0.4983	0.000382	0.00191	161.90	9.52	9.97
40,000	871.62	871.43	871.525	−0.250	0.4983	0.000502	0·00201	164.58	12.20	12.46
30,000	871.57	871.38	871.475	−0.300	0.4982	0.000602	0.00181	168.23	15.86	16.61
20,000	871.42	871.22	871.32	−0.460	0.4981	0.000924	0.00185	176.22	23.85	24.90
10,000	871.01	870.82	870.915	−0.865	0.4980	0.00174	0.00174	200.00	47.63	49.80
∞	871.89	871.67	871.78		0.4984			152.37		
						Mean capacity	0.00186			

succession in series with the condenser, the latter having a capacity of 2 microfarads. The changes in r were determined as before, and the changes in L_0 resulting therefrom were computed, and these compared with $p^2 r_1 C^2 L$ (formula 19). The changes in Q were computed from equation (18) and compared with the changes in the noninductive part of the arm Q, namely, Q'.

These measurements were made before the introduction into the bridge of the slide wire, which permits settings to 0.001 ohm, and hence were not quite as accurate as those previously given. Nevertheless the agreement between the observed and calculated values is excellent.

Table VII gives the results of measurements made to verify formulæ 22 and 23. Two series of measurements were made. In one a megohm box, consisting of 10 coils of 100,000 ohms each, was used to give resistances varying from 1,000,000 to 12,500 ohms, using various combinations of coils in series and in parallel. In the other a box of 10 coils of 10,000 ohms each was used, the coils being used in series only. In both cases the capacity of the coils produces a distinct effect on the measured value of L_0, and hence we have assumed formula (23) to be correct and have calculated the values of the capacities of the coils, which will account for the differences observed. They are given in the seventh column, and the capacity of one coil deduced from the measured capacity of the several coils is given in the eighth column. These deduced capacities per coil, averaging 0.00186 microfarad in the case of the second box, agrees quite well with an independent measurement made some months ago by a dynamometer method which gave 0.00195 microfarad for the average. It will be noticed that the changes in Q are considerable and that the observed and calculated values agree quite closely. These differences, however, can not be determined with great accuracy, as the resistance of the inductive coil (wound with copper wire) varies from time to time; hence Q' varies from this cause when Q is constant. The results, however, abundantly justify the formulæ (22 and 23).

10. GRAPHICAL SOLUTION OF THE ANDERSON BRIDGE.

A graphical solution of Anderson's bridge is interesting and shows very simply some of the results derived above analytically. If the bridge is balanced and no current is flowing through the galvanometer, we may consider that the connection $E D$ is removed. The condition of the bridge is (the same electromotive force acting on the upper half of the bridge from A to B as on the lower half) that E and D are always at the same potential.

To construct the electromotive-force diagram for the upper half of the bridge, we lay off $C B$ (fig. 11) to represent the emf. acting on the

arm R of the bridge. The same electromotive force acts on the branches $C\,E\,B$, consisting of the resistance r and the condenser of impedance $\dfrac{1}{pC}$. Therefore, in a semicircle described on $C\,B$ as a diameter construct a right-angled triangle $C\,E\,B$, the sides $C\,E$ and $B\,E$ being proportional to r and $\dfrac{1}{pC}$, respectively. $C\,E$ is then equal to e_5, the emf. acting on r, the fifth arm of the bridge, and $E\,B = e_6$,
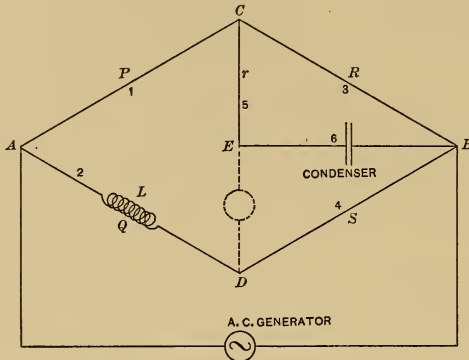
FIG. 10.—Anderson Bridge. When balanced, galvanometer may be removed.

the emf. on the condenser. $C\,c'$ and $C\,b$ are the currents in r and R, respectively (calculated from the electromotive forces and resistances), and their resultant $C\,b'$ equals the current in the arm P—that is, i_1. Of course i_5 is also the current through the condenser. The emf. acting on the arm P is in phase with the current, $C\,b'$, since P is non-inductive. Therefore, if we project $C\,b'$ backward to A, so that
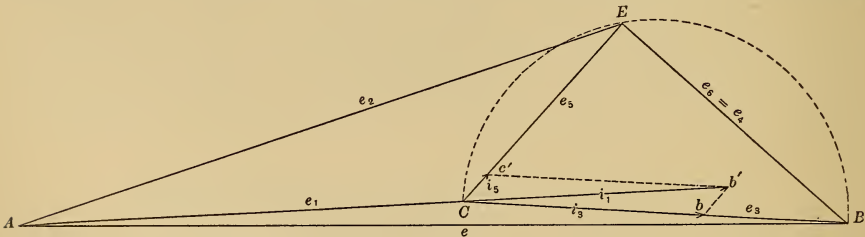
FIG. 11.—Vector diagram of Anderson Bridge.

$C\,A = i_1 \times P$, $C\,A$ is the emf. e_1 acting on P, and the vector sum of $A\,C$ and $C\,B$, or $A\,B$, will be the total emf. on the bridge.

Since the lower half of the bridge has the same emf. acting on it and the point D has the same potential as E, it is evident that the triangle $A\,E\,B$ is also the emf. triangle of the lower half $A\,D\,B$. This enables us to find graphically the inductance L in the branch Q.

Lay off DB (fig. 12) equal to EB (fig. 11), since the same emf. acts on these two branches, which terminate at the common point B, their initial ends E and D having the same potential. Draw lines DA and BA so that the triangle ADB is equal to the triangle AEB of fig. 11. ADB is then the emf. triangle of the lower half of the bridge, and AD is the emf. e_2 expended on the branch Q. Of this, DG, perpendicular to DB, overcomes the reactance pL, and AG, perpendicular to DG, overcomes the resistance Q. This construction gives L when p is known.

But p need not necessarily be known, as the values of the capacity and resistances of the bridge are independent of p. The distribution
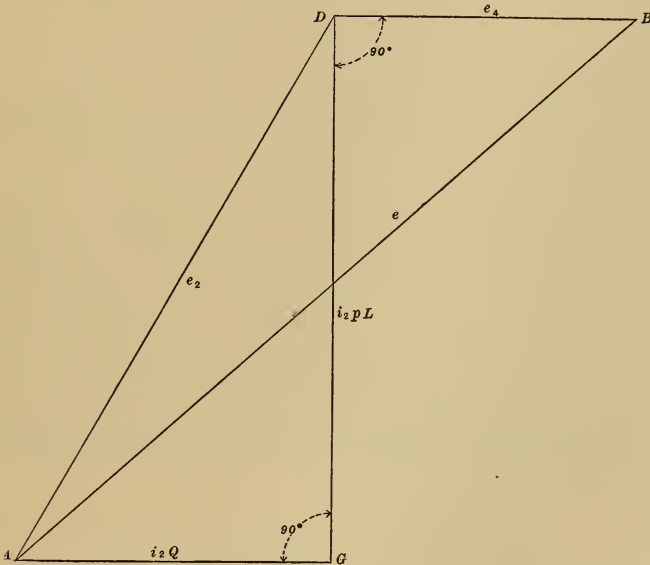
FIG. 12.—Graphical solution of Anderson Bridge.

of electromotive forces is, however, affected by the frequency, and hence the emf. triangle depends on p. If, however, any convenient value of p be assumed in constructing figures 11 and 12, the same value of L will be derived from DG; that is, DG will always come out proportional to p.

The inductance L, derived from DG, is the total inductance of the branch Q; hence, if that part of the resistance of Q not included in the inductance coil to be measured possesses positive or negative inductance (l_2), this must be subtracted from the measured value (L_o) to obtain the true inductance of the coil L.

If the branch S contains inductance l_4, DBB' will be its voltage

triangle, and the angle ϕ_4 will be $\dfrac{pl_4}{S}$.   The triangle ADB, which still represents the distribution of voltages, both in the upper and lower
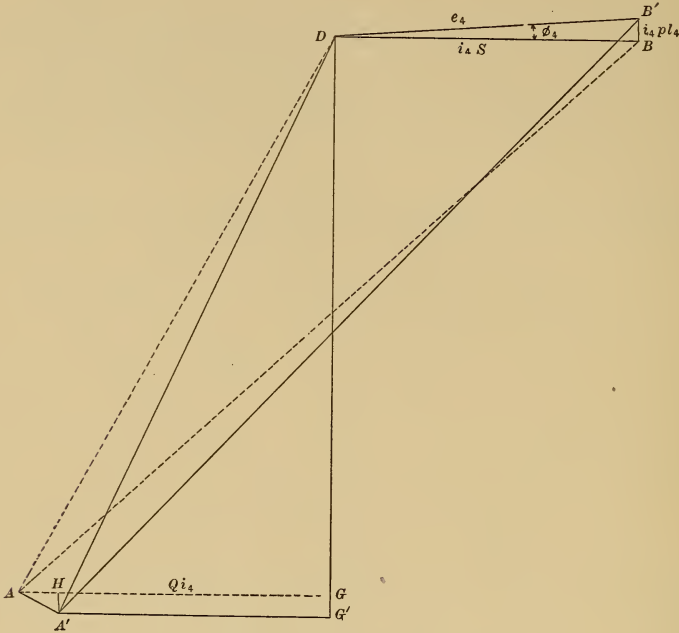
Fig. 13.—Solution when arm S has inductance l_4.

halves of the bridge, will therefore be rotated through the angle ϕ_4 into the position $A'DB'$, where

$$AA' = AD \times \phi_4$$
$$A'H = AA' \sin A'AH = AA' \sin ADG$$

Hence, $\quad GG' = AD \times \phi_4 \times \dfrac{AG}{AD} = \phi_4 \times Qi_4 = \varDelta(pL)i_4$

$\therefore \qquad \varDelta L = \dfrac{Q}{p}\phi_4 = \dfrac{l_4}{S}Q.$

This is the value of the correction due to l_4 found above analytically and given by equations (11) and (14), neglecting small quantities of the second order occurring in the β term.   A similar construction obviously applies to the case of capacity in the resistances; that is, to negative inductance.

11. RESISTANCE IN SERIES AND IN PARALLEL WITH THE CONDENSER.

Fig. 14 is the electromotive-force diagram for the Anderson bridge on which 100 volts is impressed, the frequency being such that

$p^2 = 500,000$ and $p = 707$, approximately, and fig. 15 is the correspond-
ing case, with 50 ohms inserted in series with the condenser. CF now
represents the fall in potential through r and r_1 (fig. 8); but since the
galvanometer is joined to E, between r and r_1, the triangle AEB, and
not AFB, represents the electromotive forces of the lower half of the
bridge. The values of the several electromotive forces and currents
have been accurately calculated from formula (19) and marked in the
figures. The effect of inserting r_1 in the condenser circuit is to
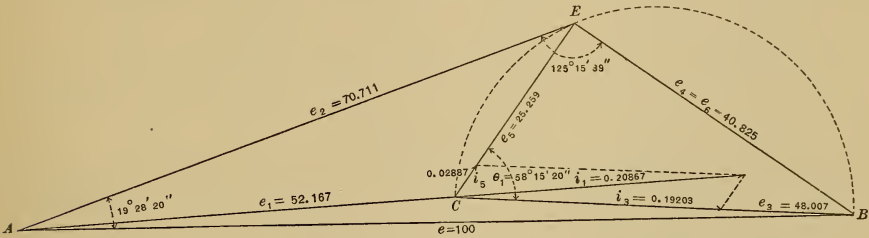
FIG. 14.—Electromotive force diagram of Anderson Bridge, having 100 volts applied to terminals.

decrease the currents i_3 and i_5. Their resultant, however, is increased,
as the angle θ_1 between them is decreased sufficiently to more than
offset the decrease in the separate currents. The current i_1 is there-
fore increased, and e_1 is increased in consequence. The side EB, the
fall in potential in S, is decreased. This shows that the current in
the lower half of the bridge is decreased, since S is noninductive.
But AE, the emf. on Q, is increased; and since the current through Q
is decreased, it follows that the impedance, and therefore the resist-
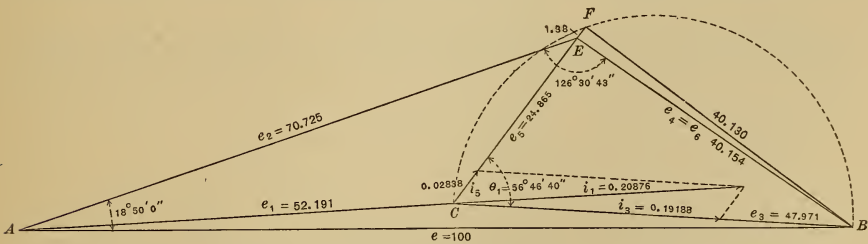
FIG. 15.—Case of Fig. 14 modified by 50 ohms resistance in series with the condenser.

ance, is increased, as is found in practice. In this case Q changes
from 500 to 525 ohms. The change in r is very slight—in this case
from 875 to 876.25. The change in L_0 is also very slight. Its value
is given by equation (19), but can not be deduced easily geometrically.
    Fig. 17 shows the effect of placing 10,000 ohms in parallel with the
condenser. In this case the current i_5 through r splits into two parts,

i_6 through the condenser and i_7 through r_2, these two components being at right angles to each other.   The result is to reduce the voltage on the condenser, and hence also the current through the condenser. The current i_5 (the sum of i_6 and i_7) is less than before, and i_3 is also
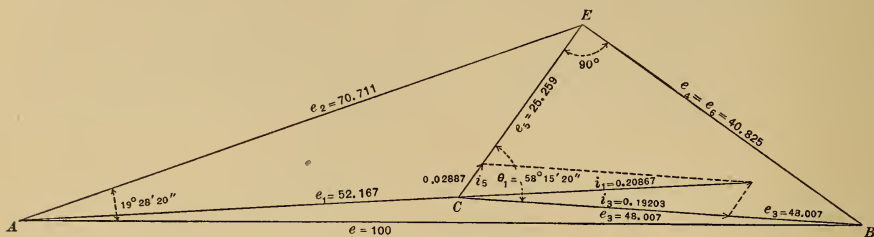
Fig. 16.—Same as Fig. 14.

less.   Nevertheless, their sum is greater, as the angle θ_1 is reduced (as in the case of resistance in series) more than enough to offset the reduction in the components.   It is remarkable, in spite of all these
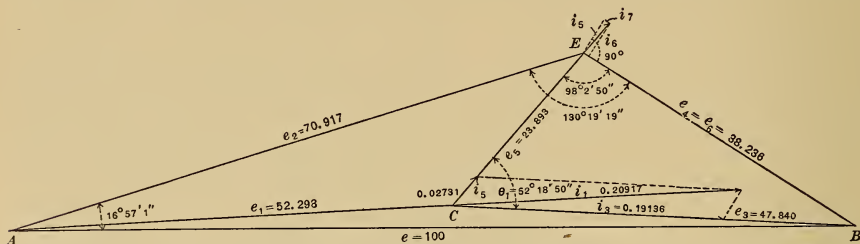
Fig. 17.—Showing effect of 10,000 ohms in parallel with the condenser.   Other conditions same as in Fig. 14.

changes and the large change in Q (in this case from 500 to 600 ohms), *that r is entirely unchanged and the observed value of L is also unchanged.*

12. MEASUREMENTS OF INDUCTANCE.

We give in Tables VIII, IX, X, and XI the results obtained on several inductance coils of 100 millihenrys and smaller, measuring them singly and in series in groups of two or three.   The two ratio coils P and R were each time reversed by a commutator and two settings of the variable resistance r made.   These two independent values of r are given in columns 3 and 4, and their mean value in column 5.   Referring to equation (13) above, if $P=R$, the inductance is

$$L= C\,S\,(2\,r+P) \qquad (24)$$

Table VIII.—Results of Measurements of Four Inductance Coils of 100 Millihenrys each. Coils Measured Singly and in Series in Groups of Two.

[$P=R=249.994$ ohms. $S=280.113$ ohms.]

No.	Coils.	r position 1.	r position 2.	r mean corrected.	$2r+P$	Temp. of condenser.	C	L_0 milli-henrys.	$L=L_0+\alpha$ milli-henrys.	L_s sums of single values.	$\Delta=L_s-L$ microhenrys.
1	A	52.638	52.598	52.628	355.250	18.7	1.00518	100.025	100.024	......	......
2	B	52.578	52.543	52.570	355.135	18.75	1.00518	99.993	99.985	......	......
3	H	52.511	52.476	52.504	355.001	18.75	1.00518	99.955	99.954	......	......
4	C	57.021	56.987	57.014	364.022	18.8	1.00517	102.494	102.486	......	......
5	C+B	234.605	234.529	234.567	719.128	18.85	1.00516	202.477	202.472	202.471	−1
6	C+A	234.673	234.604	234.638	719.271	18.85	1.00516	202.504	202.504	202.510	+6
7	H+A	230.151	230.082	230.116	710.227	18.9	1.00515	199.969	199.971	199.978	+7
8	H+B	230.085	230.015	230.050	710.094	18.9	1.00515	199.932	199.930	199.939	+9
											mean = +5
9	A	52.651	52.613	52.642	355.278	18.9	1.00515	100.031	100.023	......	......
10	H	52.507	52.471	52.499	354.992	18.95	1.00514	99.950	99.949	......	......
11	B	52.573	52.534	52.564	355.121	18.95	1.00514	99.986	99.983	......	......
12	C	57.027	56.991	57.019	364.032	19.0	1.00514	102.494	102.485	......	......
13	C+H	234.526	234.465	234.500	718.995	19.0	1.00514	202.435	202.434	202.434	0
14	C+A	234.680	234.608	234.644	719.282	19.0	1.00514	202.516	202.508	202.508	0
15	B+A	230.219	230.146	230.182	710.359	19.0	1.00514	200.004	199.999	200.006	+7
16	B+H	230.083	230.015	230.049	710.092	19.0	1.00514	199.928	199.927	199.932	+5
											mean +3

Table IX.—Results of Measurements on a Wide Range of Coils to Test the Bridge.

No.	Coils.	r position 1.	r position 2.	r mean corrected.	$2r+P$	$P=R$	S and C	L_0 milli-henrys.	$L=L_0+\alpha$ milli-henrys.	L_s sums of single values.	$\Delta=L_s-L$ micro-henrys.
1	0.1 A	66.294	66.256	66.288	382.570	250.007	260.216	100.040	100.040		
2	0.1 H	66.148	66.108	66.142	382.277		1.00492 mf	99.964	99.964		
3	A+H	257.447	257.368	257.416	764.825			199.999	199.998	200.004	+6
4	0.1 A	66.301	66.266	66.297	382.591		260.216	100.046	100.039		
5	0.1 B	66.220	66.184	66.215	382.423	250.007	1.00492 mf	100.002	99.998		
6	A+B	257.523	257.437	257.488	764.969			200.036	200.034	200.037	+3
7	50 A	11.732	11.706	11.729	273.451		260.216	50.126	50.119		
8	50 B	11.465	11.442	11.464	272.921	250.007	0.70445 mf	50.029	50.021		
9	50 A+50 B	148.191	148.138	148.174	546.341			100.149	100.144	100.140	−4
10	10 M. H. A	0.224	0.216	0.229	100.481		99.084	10.0050	10.0052		
11	10 M. H. B	0.232	0.221	0.236	100.495	100.023	1.00492 mf	10.0064	10.0066		
12	A+B	50.482	50.458	50.483	200.989			20.0128	20.0136	20.0118	−1.8
13	5 M. H. A	12.928	12.913	12.931	125.888		99.084	5.0026	5.0029		
14	5 M. H. B	13.016	13.006	13.022	126.067	100.023	0.40106 mf	5.0097	5.0037		
15	A+B	75.847	75.821	75.847	251.717			10.0029	10.0045	10.0066	+2.1
16	1 M. H. A	0.405	0.395	0.410	100.833		99.084	1.0033	1.0011		
17	1 M. H. B	0.520	0.515	0.528	101.079	100.023	0.10042 mf	1.0057	1.0036		
18	A+B	50.595	50.570	50.595	201.213			2.0022	2.0016	2.0047	+3.1
19	0.5 M. H. A	0.575	0.565	0.570	51.17		49.04	0.5041	0.5048		
20	0.5 M. H. B	0.480	0.475	0.478	50.98	50.01	0.20089 mf	0.5023	0.5025		
21	A+B	26.03	26.02	26.025	102.08			1.0057	1.0050	1.0073	+2.3

The values of $2r+P$ are given in column 6 of the tables. The temperature of the condenser is given in column 7 of Table VIII, and its capacity at the given temperature in column 8. In Tables IX, X, and XI the temperature of the condenser was kept constant. The values of L, calculated from equation (24), are given in column 10, and in column 11 are given the corrected values of L, obtained by applying the value of α, which is sometimes positive and sometimes negative. In column 12 the sums (L_s) of the separate values of L are given, for comparison with the values opposite to them in column 11, obtained by measuring the coils in groups of two or three. The differences between the values of L_s and L are given in column 13. These differences are discrepancies which indicate errors in the measured values of the inductances, but which amount on the average to only five parts in 200,000 for the 100-millihenry coils, and are sometimes positive and sometimes negative. They represent the errors of observation, the errors of the resistances, and the imperfections in the corrections for inductance and capacity of the arms of the bridge. The errors due to differences in P and R are eliminated by the commutator; the error due to inductance in r we have seen to be inappreciable; the condenser is so good that no correction needs to be applied for its very slight absorption and leakage. In the second half of the table the results are given for a remeasurement of the coils under conditions similar to those under which the first measurements were made. It will be noticed that the values agree very closely.

In Table IX the results are given of measurements on coils of 100, 50, 10, 5, 1.0, and 0.5 millihenrys, using 4 of 100 millihenrys and 2 each of the other denominations, measuring them separately and in series in each. The residual differences between the corresponding values in columns 11 and 12, and which amount to a very few microhenrys, are as small as could be expected.

In Table X the results are given of measurements on a coil C of nominally 100 millihenrys, and of two coils, A and B, of 50 millihenrys each, measured together as one coil and in series with the coil C. The residual differences between the sums of the separate values (column 12) and the measured sum (column 11) are given in column 13. They average about 4 in 200,000, being sometimes positive and sometimes negative. In Table XI similar results are given for the two 50-millihenry coils measured separately and in series, with the residuals in column 13 averaging 6 parts in 100,000.

In both these tables the measurements are made with three different values of the resistances of the bridge (P, R, and S being equal in every case) and three different currents for each separate value of the

Table X.—Results of Measurements on a Coil of 100 Millihenrys and on Two 50 Millihenry Coils in Series.

1	2	3	4	5	6	8	9	10	11	12	13
No.	Coils.	r position 1.	r position 2.	r mean corrected.	$2r+P$	P, S and C	Total bridge current amperes.	L_0 millihenrys.	$L=L_0+\alpha$ millihenrys.	L_s sums of single values.	$\Delta=L_s-L$ microhenrys.
1	C+(A+B)	278.080	278.050	278.098	806.214			202.607	202.605	202.597	−8
2	(A+B)	74.193	74.154	74.191	398.400	P=R=250.018	0.12	100.124	100.120		
3	C	78.892	78.850	78.888	407.794	S=250.233		102.485	102.477		
4	C+(A+B)	278.109	278.027	278.095	806.208			202.614	202.612	202.606	−6
5	(A+B)	74.204	74.164	74.201	398.420	C=1.00433	0.20	100.130	100.126		
6	C	78.891	78.862	78.893	407.804			102.488	102.480		
7	C	78.883	78.856	78.887	407.792			102.485	102.476		
8	(A+B)	74.177	74.162	74.187	398.392		0.30	100.122	100.118		
9	C+(A+B)	278.064	278.035	278.075	806.168			202.604	202.602	202.594	−6
10	C	154.872	154.831	154.868	509.762			102.493	102.469		
11	(A+B)	149.004	148.966	149.002	498.030	P=R=200.026	0.12	100.135	100.114		
12	C+(A+B)	403.837	403.750	403.806	1,007.638	S=200.194		202.597	202.578	202.583	+5

No.												
13	C+(A+B)	403.870	403.800	403.848	1,007.722	C=1.00433		202.614	202.595	202.597	
14	(A+B)	149.020	148.986	149.020	498.066		0.20	100.142	100.121		+2
15	C	154.881	154.856	154.885	509.796			102.500	102.476		
16	C	154.869	154.856	154.875	509.776			102.496	102.472		
17	(A+B)	148.987	148.988	149.005	498.036		0.30	100.136	100.115		
18	C+(A+B)	403.823	403.823	403.836	1,007.698			202.609	202.590	202.587	−3
19	C+(A+B)	596.680	596.750	596.743	1,343.494	P=R=150.008		202.588	202.584	202.582	−2
20	(A+B)	256.943	256.966	256.982	663.972		0.30	100.119	100.114		
21	C	264.755	264.785	264.797	679.602	S=150.138		102.476	102.468		
22	C	264.740	264.760	264.777	679.562			102.470	102.462		
23	C+(A+B)	596.680	596.677	596.706	1,343.420		0.20	202.572	202.569	202.573	+4
24	(A+B)	256.950	256.934	256.971	663.950			100.116	100.111		
25	(A+B)	256.893	256.875	256.913	663.834			100.098	100.094		
26	C+(A+B)	596.600	596.545	596.600	1,343.208		0.12	202.540	202.537	202.543	+6
27	C	264.725	264.697	264.738	679.484			102.458	102.449		

Table XI.—Results of Measurements on Two Inductance Coils of 50 Millihenrys Each, Coils Measured Singly and in Series.

[Temperature of condenser = 24.4°.]

1	2	3	4	5	6	8	9	10	11	12	13
No.	Coils.	r position 1.	r position 2.	r mean corrected.	$2r+P$	P and S and C	Total bridge current amperes.	L_0 millihenrys.	$L=L_0+\alpha$ millihenrys.	L_8 sums of single values millihenrys.	$\Delta=L_8-L$ microhenry.
1	50A	201.705	201.766	201.765	503.546	....	....	50.113	50.111	....	....
2	50B	201.208	201.273	201.267	502.550	P=100.016	0.30	50.014	50.010	....	....
3	50A+50B	452.915	453.037	452.988	1,005.992	S= 99.091	....	100.111	100.111	100.121	+10
4	50A+50B	452.885	452.941	452.923	1,005.862	R=100.016	....	100.106	100.104	....	....
5	50B	201.195	201.235	201.245	502.506	....	0.20	50.010	50.006	....	....
6	50A	201.675	201.725	201.732	503.480	C=1.00433	....	50.106	50.104	50.110	+ 6
7	50A	201.645	201.645	201.675	503.366	....	....	50.095	50.092	....	....
8	50B	201.145	201.155	201.180	502.376	....	0.12	49.996	49.993	....	....
9	50A+50B	452.740	452.790	452.775	1,005.566	....	....	100.074	100.074	100.085	+11
10	50A	91.133	91.110	91.143	332.294	P=R=150.008	....	50.106	50.100	....	....
11	50B	90.808	90.792	90.821	331.650	....	0.12	50.009	50.001	....	....
12	50A+50B	256.910	256.870	256.917	663.842	S=150.138	....	100.100	100.096	100.102	+ 6

No.	Coil										
13	50A+50B	256.963	256.933	256.975	663.958			100.117	100.112	100.112	0
14	50B	90.820	90.817	90.839	331.686	C=1.00433	0.20	50.014	50.006		
15	50A	91.148	91.140	91.165	332.338			50.113	50.106		
16	50A	91.144	91.148	91.167	332.342			50.113	50.107		
17	50B	90.824	90.822	90.844	331.696		0.30	50.016	50.008		
18	50A+50B	256.939	256.962	256.977	663.962			100.118	100.113	100.115	+2
19	50A+50B	149.003	148.989	149.013	498.052	P=R=200.026		100.139	100.120	100.112	−8
20	50B	24.391	24.389	24.403	248.832		0.30	50.030	50.007		
21	50A	24.636	24.628	24.645	249.316	S=200.194		50.128	50.105		
22	50A	24.637	24.629	24.646	249.318			50.128	50.105		
23	50B	24.396	24.385	24.403	248.832	C=1.00433	0.20	50.030	50.007		
24	50A+50B	149.005	148.986	149.012	498.050			100.138	100.119	100.112	−7
25	50A+50B	148.985	148.961	148.990	498.006			100.130	100.111	100.114	+3
26	50B	24.402	24.385	24.406	248.832		0.12	50.030	50.007		
27	50A	24.647	24.630	24.651	249.328			50.130	50.107		

Table XII.—Summary of Values of Inductances Shown in Table X, with the Deviations from the Mean.

1	2	3	4	5	6	7	8	9
No.	$P=R=S$.	Total current amperes	$(A+B)$ milli-henrys.	Deviation from mean.	C milli-henrys.	Deviation from mean.	$C+(A+B)$ milli-henrys.	Deviation from mean.
1	250	0. 12	100. 120	0. 005	102. 477	0. 007	202. 605	0. 019
2	"	0. 20	100. 126	. 011	102. 480	. 010	202. 612	. 026
3	"	0. 30	100. 118	. 003	102. 476	. 006	202. 602	. 016
4	200	0. 12	100. 114	. 001	102. 469	. 001	202. 578	. 008
5	"	0. 20	100. 121	. 006	102. 476	. 006	202. 595	. 009
6	"	0. 30	100. 115	0	102. 472	. 002	202. 590	. 014
7	150	0. 12	100. 094	. 021	102. 449	. 021	202. 537	. 049
8	"	0. 20	100. 111	. 004	102. 462	. 008	202. 569	. 017
9	"	0. 30	100. 114	. 001	102. 468	. 002	202. 584	. 002
		Means.	100. 115	. 006	102. 470	. 007	202. 586	. 018

Table XIII.—Summary of Values of Inductances Shown in Table XI, with the Deviations from the Mean.

1	2	3	4	5	6	7	8	9
No.	$P=R=S$.	Total current am-peres.	A milli-henrys.	Deviation from mean.	B milli-henrys.	Deviation from mean.	$A+B$.	Deviation from mean.
1	100	0. 12	50. 092	0. 012	49. 993	0. 012	100. 074	0. 033
2	"	0. 20	50. 104	. 000	50. 006	. 001	100. 104	. 003
3	"	0. 30	50. 111	. 007	50. 010	. 005	100. 111	. 004
4	150	0. 12	50. 100	. 004	50. 001	. 004	100. 096	. 011
5	"	0. 20	50. 106	. 002	50. 006	. 001	100. 112	. 005
6	"	0. 30	50. 107	. 003	50. 008	. 003	100. 113	. 006
7	200	0. 12	50. 107	. 003	50. 007	. 002	100. 111	. 004
8	"	0. 20	50. 105	. 001	50. 007	. 002	100. 119	. 012
9	"	0. 30	50. 105	. 001	50. 007	. 002	100. 120	. 013
		Means.	50. 104	. 004	50. 005	. 004	100. 107	. 010

resistances.   This gives nine measurements of each coil and of the sum of the coils.   These nine values of each coil and of their sums are given in Tables XII and XIII, with their deviations from the mean. In most cases the values are a little smaller with the smaller currents and smaller resistances.   We have not yet ascertained why this is so; perhaps the resistances were not as accurately known as we supposed.

In Table XIV the measurements of April 21 are given, three coils of 1 henry each being taken singly and in series in groups of two and three.

The separate values found during the day for the three coils are given in Table XV.   The small increase in the value of L may be due in part to uncertainty in the change of capacity of the condenser. The latter changed in temperature, according to the thermometer, by $0°.75$, and that corresponds to 11 parts in 100,000 in the capacity.   The slight progressive changes in the values of the inductances of the coils may be accounted for by a quarter of a degree greater change in the temperature of the condenser than indicated by the thermometer, or a slightly greater temperature coefficient.   We shall investigate this further, keeping the temperature of the condenser constant, to decide whether the coils really change in the manner indicated.

Table XV.—Results of the Determinations of the Inductance of Three Coils of 1 Henry each, April 21, 1905.

Coil F.	Coil S.	Coil C.
Henrys.	*Henrys.*	*Henrys.*
0. 99890	0. 99969	1. 01420
. 99892	. 999715	1. 01421
. 99894	. 99970	1. 01422
. 99895	. 999715	1. 01422
. 99895	. 99972	

The regularity of this progressive change shows that some common cause affects all the measurements, but the changes are very small indeed, amounting to only a few parts in a hundred thousand.   The sensitiveness of the bridge is well shown by these results, and if we can eliminate the small residual errors of the bridge completely, it will make it possible to measure inductances with far greater accuracy than has been done heretofore.

In order to make these measurements under the most favorable circumstances, we have designed and are now constructing a bridge espe-

Table **XIV.**—Results of Measurements on Three Inductance Coils of 1 Henry each. Coils Measured Singly and in Series in Groups of Two and Three, April 21, 1905.

$[P=R=799.819$ ohms. $S=800.425$ ohms.$]$

1	2	3	4	5	6	8	9	11	12	13
No.	Coils.	r position 1.	r position 2.	r mean corrected.	$2r+P$	Temperature of condenser.	C	L measured value.	L_s sum of single values.	$\Delta=L_s-L$
		Ohms.	Ohms.	Ohms.	Ohms.	°	Microfarads.	Henrys.	Henrys.	Henrys.
1	F	220.953	220.963	220.960	1241.740	19.85	1.00501	0.99890		
2	C	230.472	230.470	230.474	1260.766	19.9	1.005005	1.01420		
3	F+C	851.365	851.375	851.343	2502.505	19.95	1.00500	2.01308	2.01310	+0.00002
4	S	221.457	221.458	221.460	1242.739	19.95	1.00500	0.99969		
5	F	220.973	220.979	220.978	1241.776	19.95	1.00500	0.99892		
6	S+F	842.310	842.295	842.275	2484.369		1.00500	1.99849	1.99861	+0.00012
7	S	221.503	221.498	221.503	1242.825	20.25	1.00495	0.999715		
8	C	230.517	230.513	230.518	1260.854	20.3	1.004945	1.01421		
9	S+C	851.930	851.930	851.903	2503.625	20.3	1.004945	2.01387	2.01393	+0.00006
10	F	221.030	221.028	221.032	1241.882	20.35	1.00494	0.99894		
11	F+S+C	1472.833	1472.823	1472.760	3745.339	20.4	1.00493	3.01264	3.01287	+0.00023
12	F+C	851.502	851.488	851.468	2502.755	20.45	1.00492	2.01313	2.01315	+0.00002

13	S	221.512	221.517	221.517	1242.853	20.5	1.004915	0.99970	----	----
14	F	221.051	221.048	221.052	1241.923	20.5	1.004915	0.99895	----	----
15	S+F	842.438	842.428	842.406	2484.631	20.5	1.004915	1.99852	1.99865	+0.00013
16	C+S	851.977	851.967	851.945	2503.709	20.6	1.00490	2.01387	2.01394	+0.00007
17	C	230.541	230.544	230.545	1260.909	20.6	1.00490	1.01422	----	----
18	S	221.526	221.531	221.531	1242.881	20.6	1.00490	0.999715	----	----
19	F+C	851.537	851.530	851.506	2502.831	20.6	. 1.00490	2.01314	2.01317	+0.00003
20	F	221.056	221.058	221.060	1241.938	20.6	1.00490	0.99895	----	----
21	C	230.545	230.549	230.550	1260.918	20.6	1.00490	1.01422	----	----
22	S	221.533	221.536	221.537	1242.893	20.6	1.00490	0.99972	----	----
23	F+C+S	1472.930	1472.910	1472.852	3745.523	20.6	1.00490	3.01270	3.01289	+0.00019

cially adapted to this work. The bridge we have been using has been made up of several different resistance boxes. The resistances were all immersed in oil, the temperatures were accurately taken, they were frequently measured against our standards, and every effort made to get precise results. But with all combined in a single set, with shorter and more permanent connections between the arms of the bridge and the residual inductances and capacities eliminated more completely, we hope to improve upon the results given above, and make sure of the value of the inductance of a standard coil to at least one part in ten thousand, barring the uncertainty in the ohm itself. When this has been accomplished we hope to determine some inductance coils, the dimensions of which shall have been accurately determined, and the inductances computed therefrom, in order to obtain a check upon the value of the ohm.

The Anderson bridge, used in the manner we have described above, is admirably adapted to practical use in measuring inductances of a great range of values. Its accuracy far exceeds most of the methods that have been proposed, and it is more convenient than any method we have employed that can compare with it in accuracy.

NOTE.—Since the above was written we have discovered that the serpentine spools on which the inductances F and S of 1 henry each and several of 0.1 henry are wound are slightly magnetic, and that their permeability is larger with higher magnetizing force. Consequently the inductances of the coils are not quite constant, but vary when the voltage on the bridge or the resistances of the arms of the bridge are varied. This would cause the sum of two inductances measured separately to be a little greater than when measured together, and this accounts for the discrepancies Δ in column 13 of Table XIV, all being of the same sign, and larger when coils F and S (both being wound on serpentine spools) are taken together than when F and C or S and C are taken together, C being wound on a mahogany spool and therefore not magnetic.

In the case of coils of 0.1 henry or less the current is reduced very little when two are put in series in the bridge, and hence the resultant change in the inductance is much less.

We give in another article in this bulletin the results of tests made to determine the magnitude of the effect of this small amount of magnetic matter in the spools.